

Quark-Hadron Duality in Electron Scattering

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Goal: QCD-based understanding of structure function data at low Q^2 and W

Outline:

- Quark-hadron (*“Bloom-Gilman”*) duality
- Duality in QCD
- Local duality & truncated moments
- Conclusions

Quark-Hadron
(*“Bloom-Gilman”*)
Duality

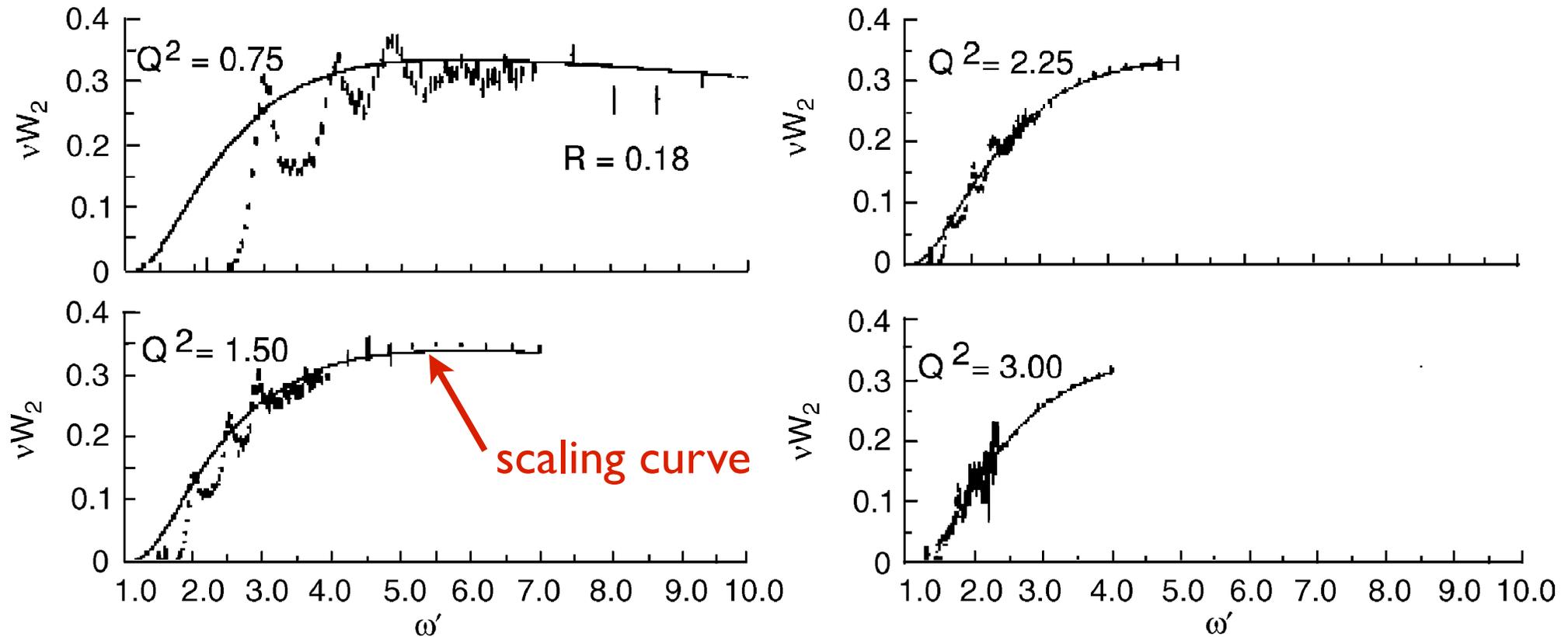
Quark-hadron duality

Complementarity between *quark* and *hadron* descriptions of observables

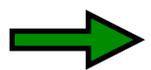
$$\sum_{\text{hadrons}} = \sum_{\text{quarks}}$$

Can use either set of complete basis states to describe all physical phenomena

Inclusive electron-proton scattering



Bloom, Gilman, Phys. Rev. Lett. 85 (1970) 1185



**resonance – scaling duality in
proton $\nu W_2 = F_2$ structure function**

Bloom-Gilman duality

Average over (strongly Q^2 dependent) resonances
 $\approx Q^2$ independent scaling function

Finite energy sum rule for eN scattering

$$\frac{2M}{Q^2} \int_0^{\nu_m} d\nu \nu W_2(\nu, Q^2) = \int_1^{\omega'_m} d\omega' \nu W_2(\omega')$$

measured structure function
(function of ν and Q^2)

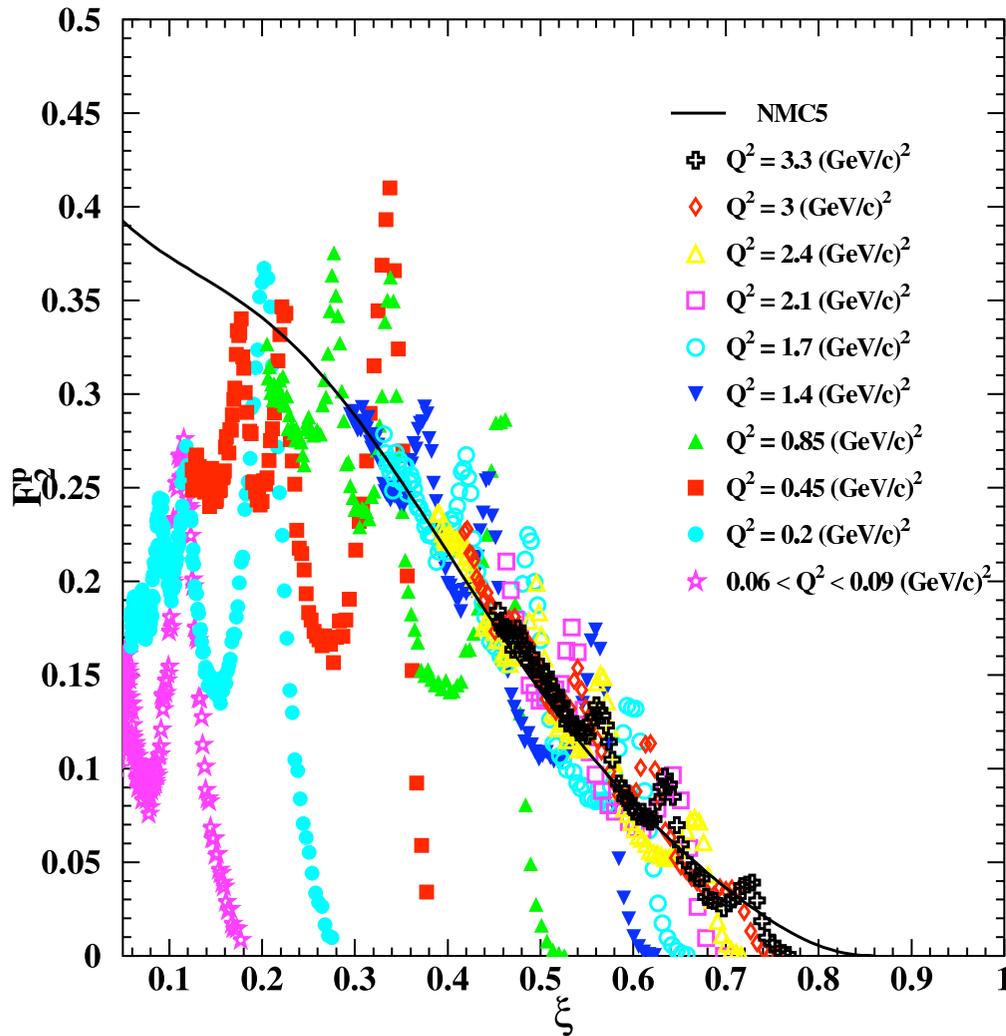
“hadrons”

$$\omega' = \frac{1}{x} + \frac{M^2}{Q^2}$$

scaling function
(function of ω' only)

“quarks”

Bloom-Gilman duality

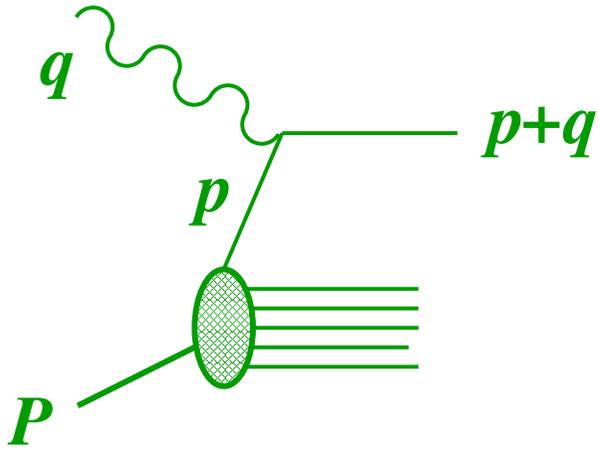


Average over
(strongly Q^2 dependent)
resonances
 \approx Q^2 independent
scaling function

Jefferson Lab (Hall C)

Niculescu et al., Phys. Rev. Lett. 85 (2000) 1182

Scaling variables



$$(p + q)^2 = m_q^2 \quad \left\{ \begin{array}{l} m_q = 0 \\ p_T = 0 \end{array} \right.$$

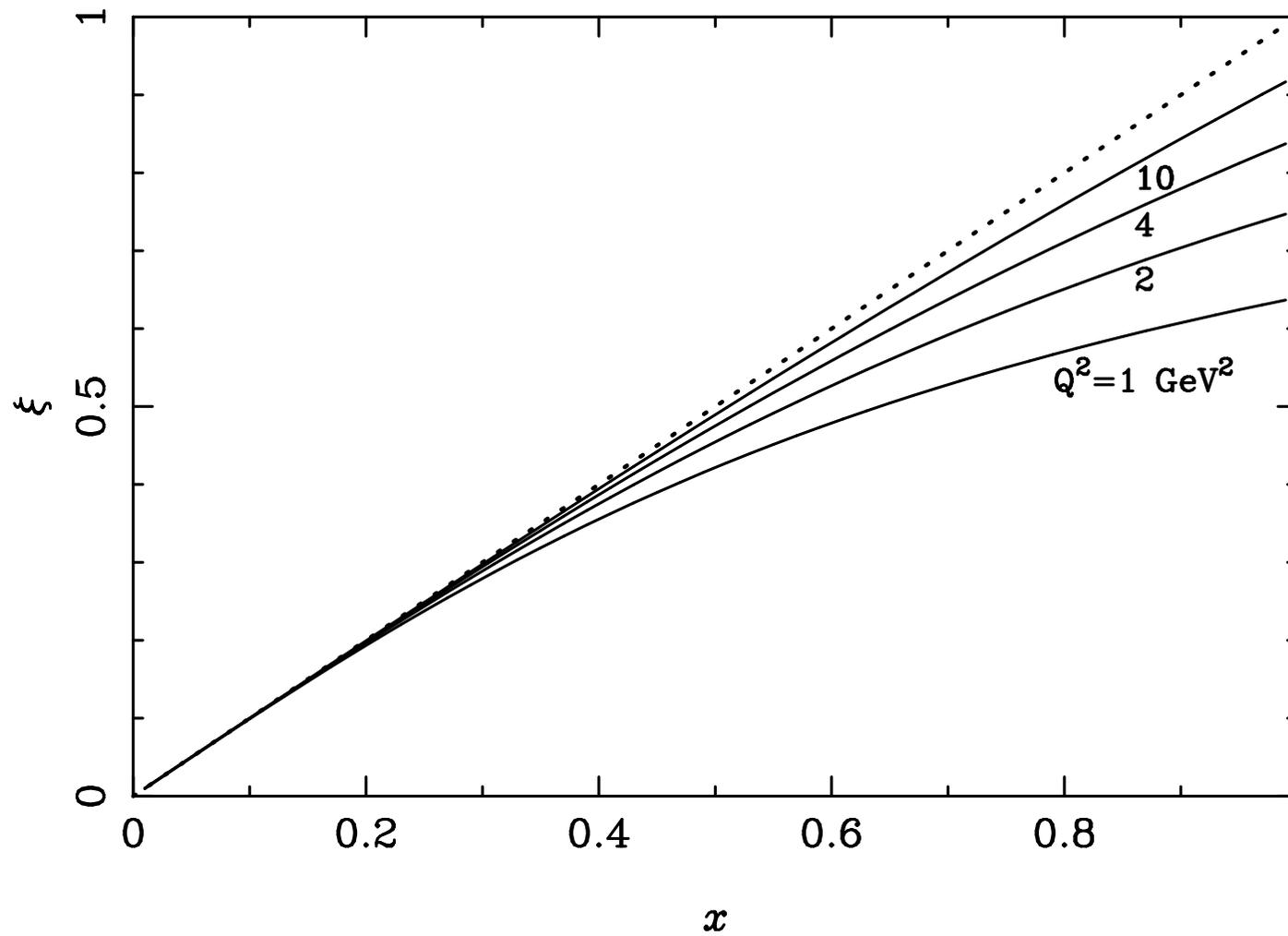
light-cone fraction of target's momentum carried by parton

$$\xi = \frac{p^+}{P^+} = \frac{p^0 + p^z}{M}$$

$$\rightarrow \xi = \frac{2x}{1 + \sqrt{1 + 4x^2 M^2 / Q^2}} \rightarrow x \text{ as } Q^2 \rightarrow \infty$$

Nachtmann scaling variable

Scaling variables



Duality in QCD

Duality in QCD

Operator product expansion

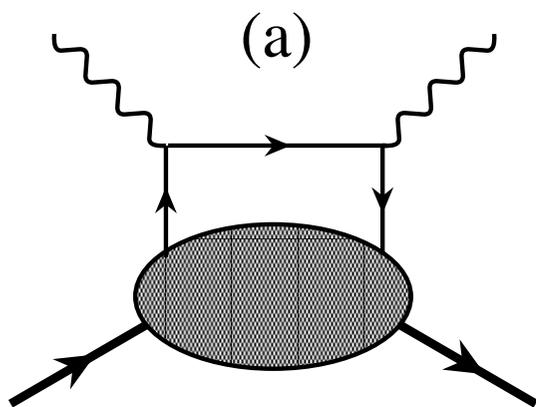
→ expand moments of structure functions
in powers of $1/Q^2$

$$\begin{aligned} M_n(Q^2) &= \int_0^1 dx x^{n-2} F_2(x, Q^2) \\ &= A_n^{(2)} + \frac{A_n^{(4)}}{Q^2} + \frac{A_n^{(6)}}{Q^4} + \dots \end{aligned}$$

matrix elements of operators with
specific “twist” τ

$\tau = \text{dimension} - \text{spin}$

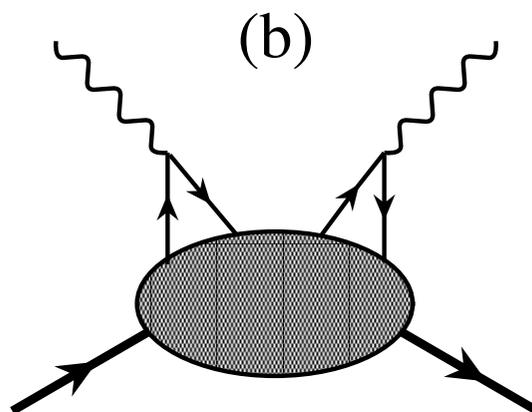
Higher twists



$$\tau = 2$$

single quark
scattering

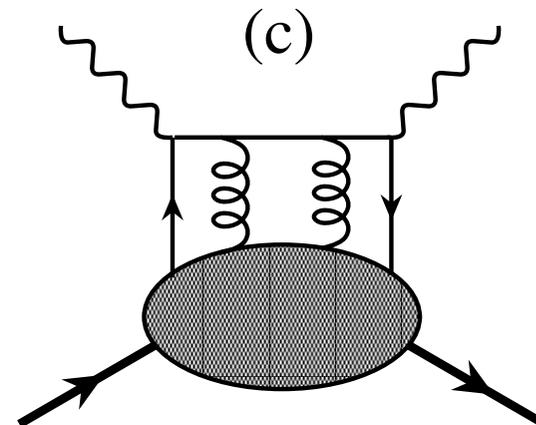
e.g. $\bar{\psi} \gamma_\mu \psi$



$$\tau > 2$$

qq and *qg*
correlations

e.g. $\bar{\psi} \gamma_\mu \psi \bar{\psi} \gamma_\nu \psi$
or $\bar{\psi} \tilde{G}_{\mu\nu} \gamma^\nu \psi$



Duality in QCD

Operator product expansion

→ expand moments of structure functions
in powers of $1/Q^2$

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If moment \approx independent of Q^2

→ higher twist terms $A_n^{(\tau>2)}$ small

Duality in QCD

Operator product expansion

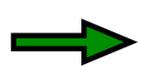
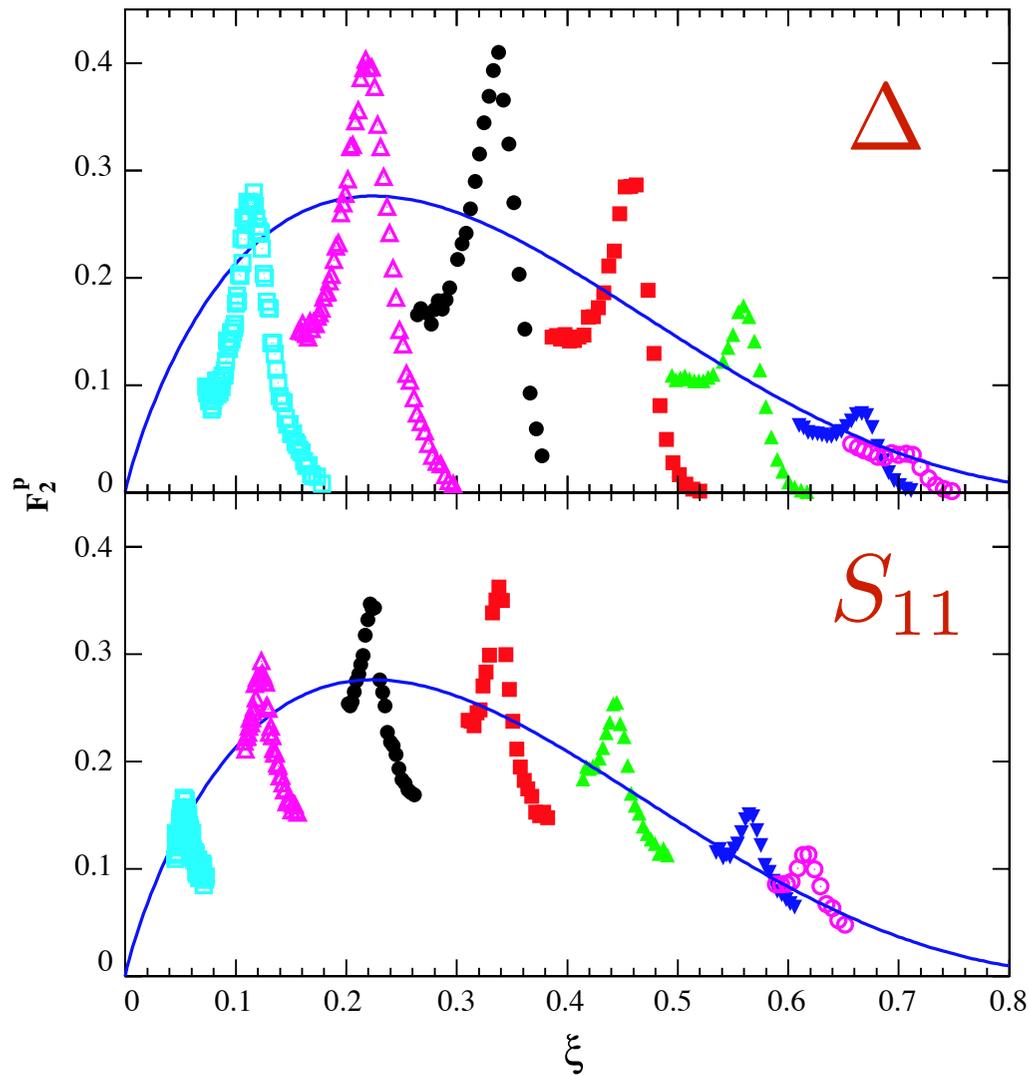
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Duality \iff suppression of higher twists

*de Rujula, Georgi, Politzer,
Ann. Phys. 103 (1975) 315*

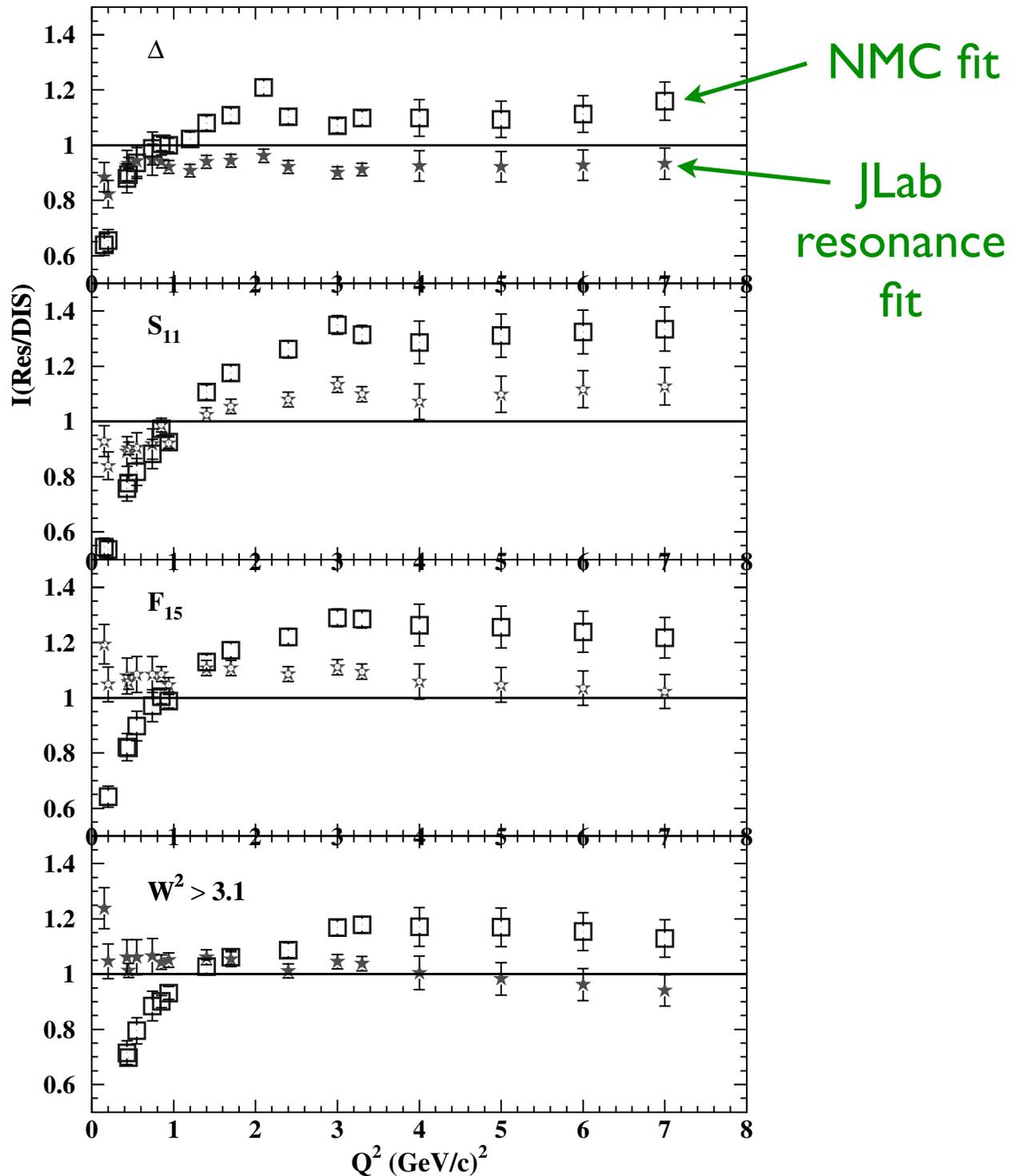
- Duality exists also in local regions, around individual resonances



“local Bloom-Gilman duality”

ratio of integrated
resonance to DIS
contributions

~10% agreement
for $Q^2 > 1 \text{ GeV}^2$



Truncated Moments

Truncated moments

- complete moments can be studied in QCD via twist expansion
 - Bloom-Gilman duality has a precise meaning
(*i.e.*, duality violation = higher twists)
- for local duality, difficult to make rigorous connection with QCD
 - *e.g.* need prescription for how to average over resonances
- truncated moments allow study of restricted regions in x (or W) within QCD in well-defined, systematic way

$$\overline{M}_n(\Delta x, Q^2) = \int_{\Delta x} dx x^{n-2} F_2(x, Q^2)$$

Truncated moments

- truncated moments obey DGLAP-like evolution equations, similar to PDFs

$$\frac{d\overline{M}_n(\Delta x, Q^2)}{d \log Q^2} = \frac{\alpha_s}{2\pi} \left(P'_{(n)} \otimes \overline{M}_n \right) (\Delta x, Q^2)$$

where modified splitting function is

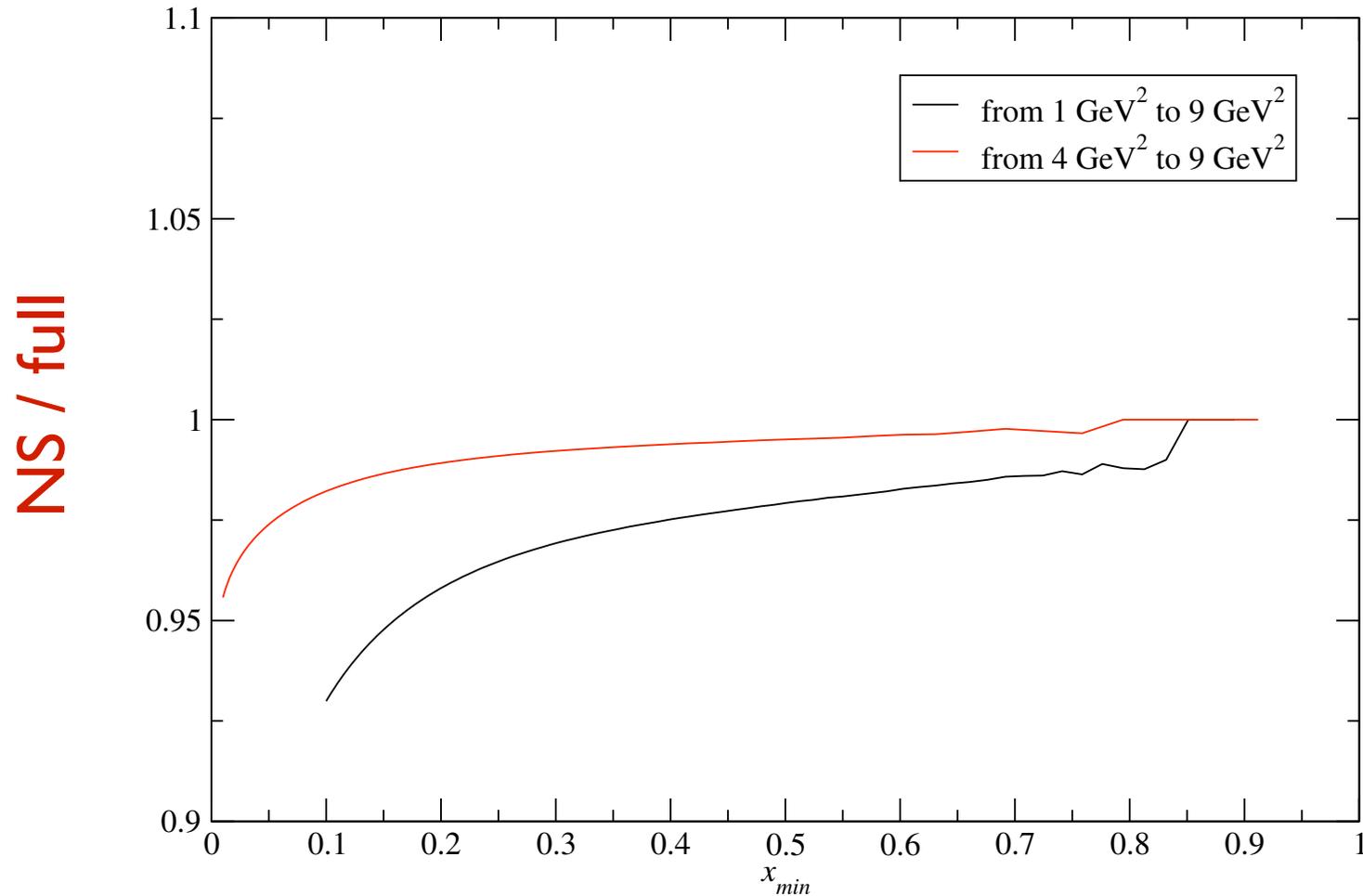
$$P'_{(n)}(z, \alpha_s) = z^n P_{NS,S}(z, \alpha_s)$$

- can follow evolution of specific resonance (region) with Q^2 in pQCD framework!
- suitable when complete moments not available

Truncated moments

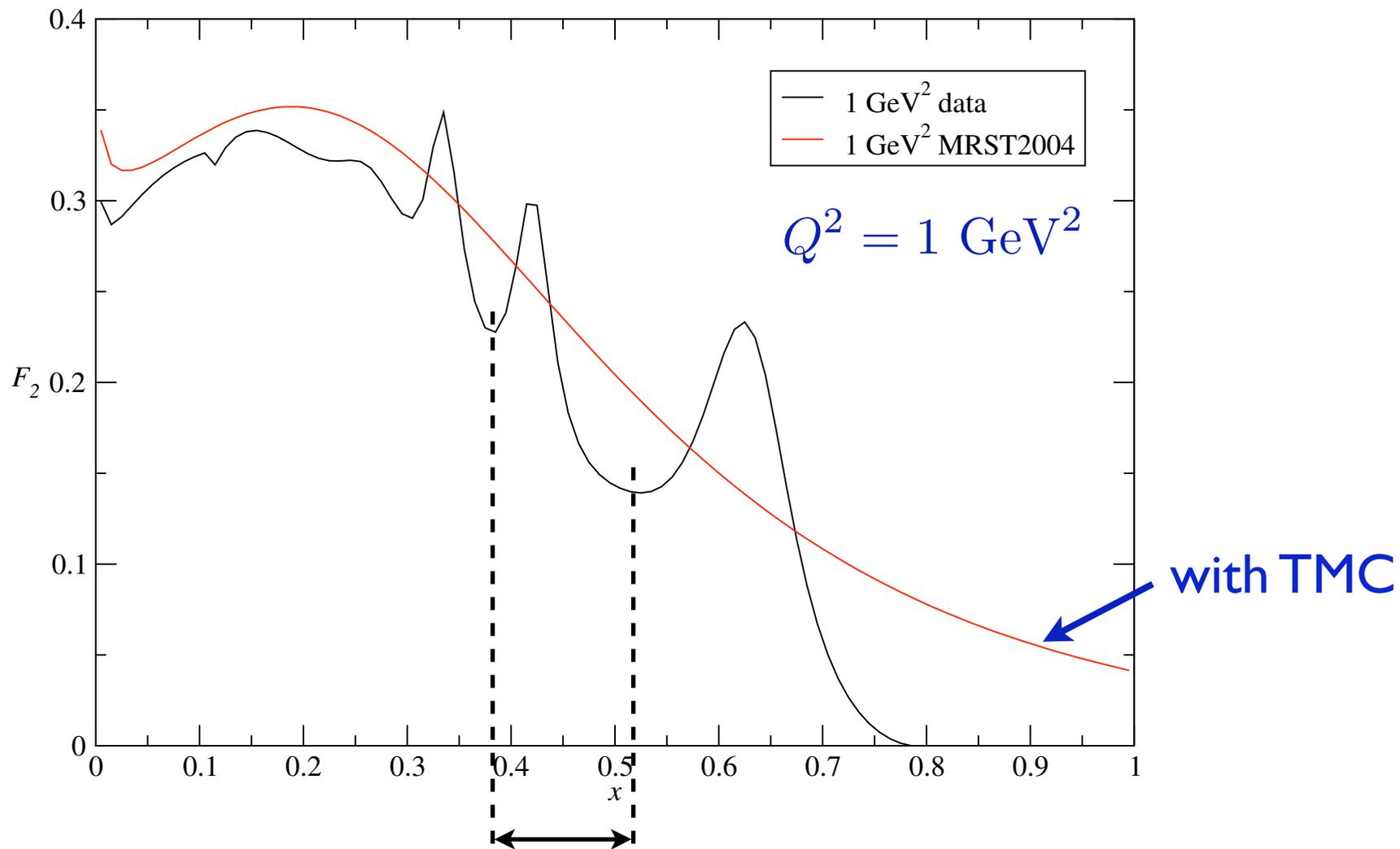
- truncated moment evolution equations exist for singlet (S) and nonsinglet (NS) separately
- for analysis of data, do not know much of experimental structure function is NS and how much is S
 - for lowest ($n=2$) truncated moment, assumption that total \approx NS is good to few % for $x_{\min} > 0.2$
 - for higher moments, small- x region is further suppressed, so that NS is a very good approximation to total

$n = 2$ truncated moment of F_2^p



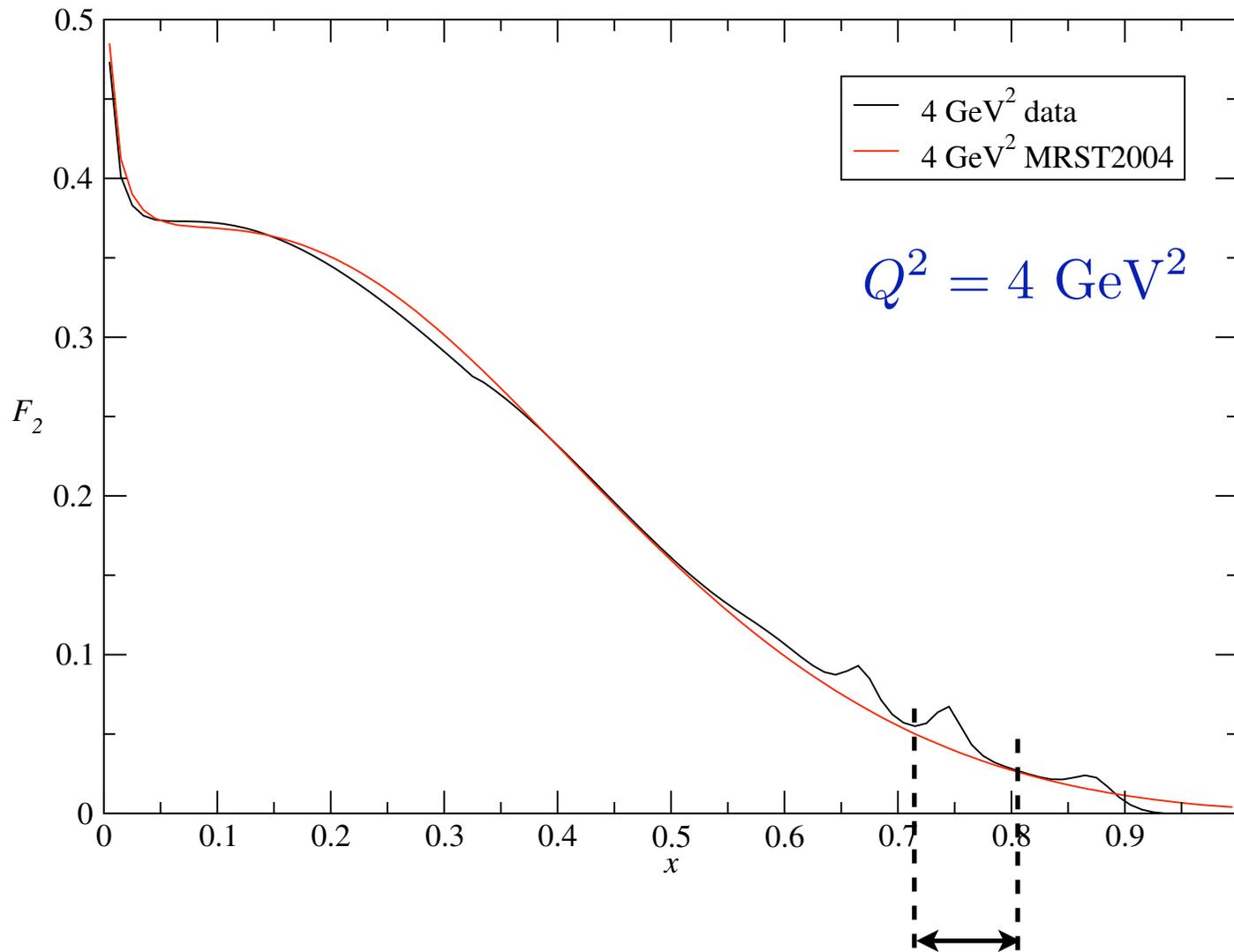
$$\Delta x = [x_{min}, 1]$$

Parameterization of F_2^p data



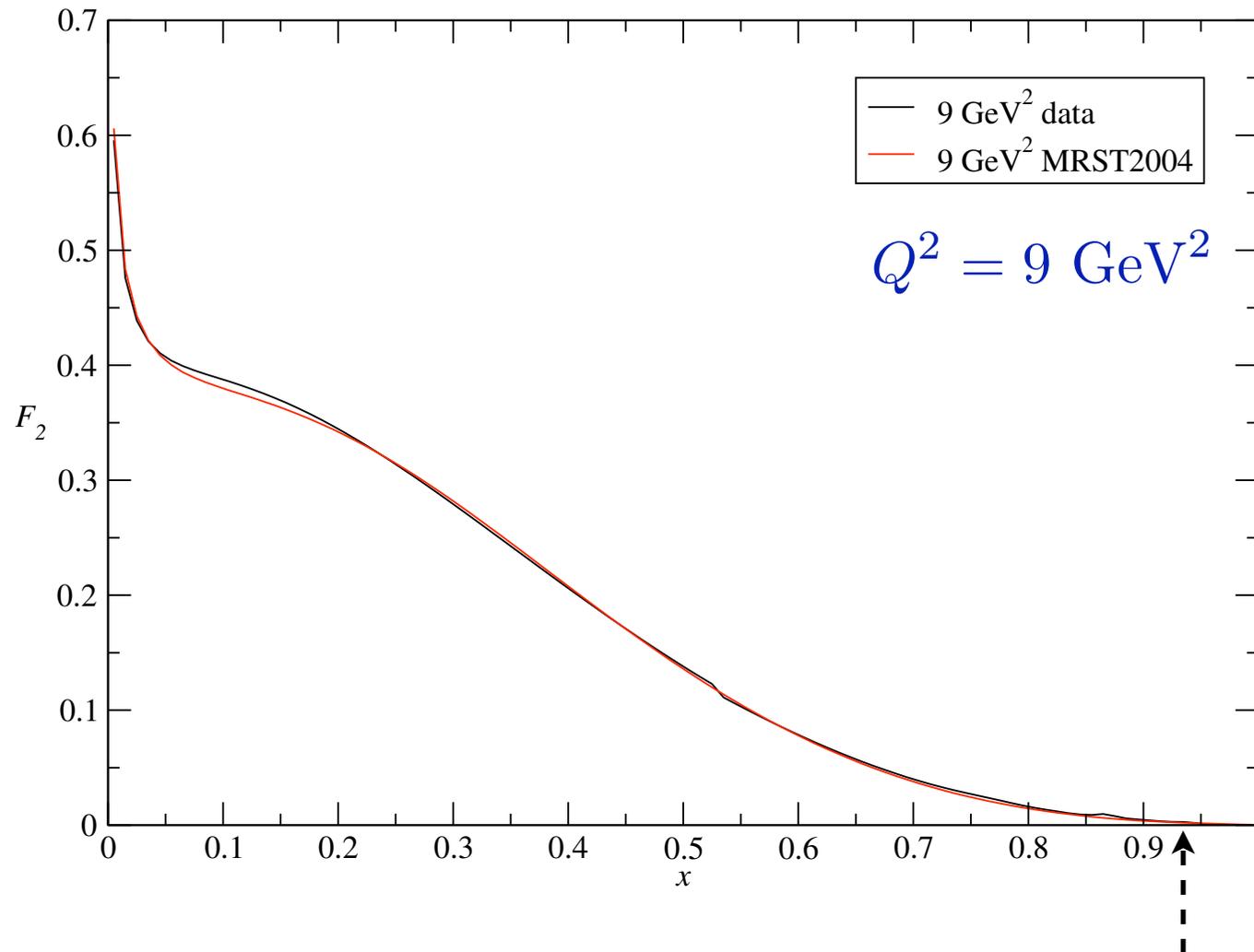
how much of this region
is leading twist ?

Parameterization of F_2^p data



how much of this region
is leading twist ?

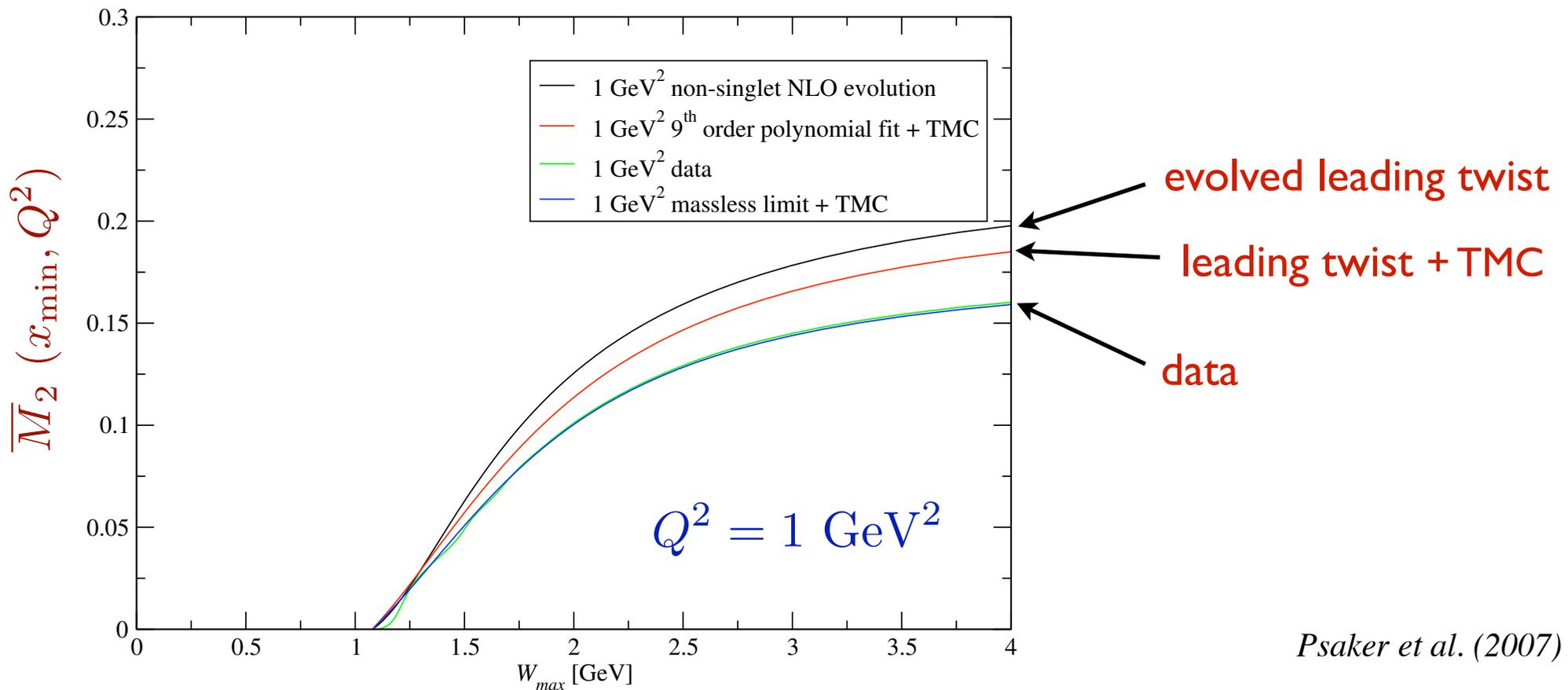
Parameterization of F_2^p data



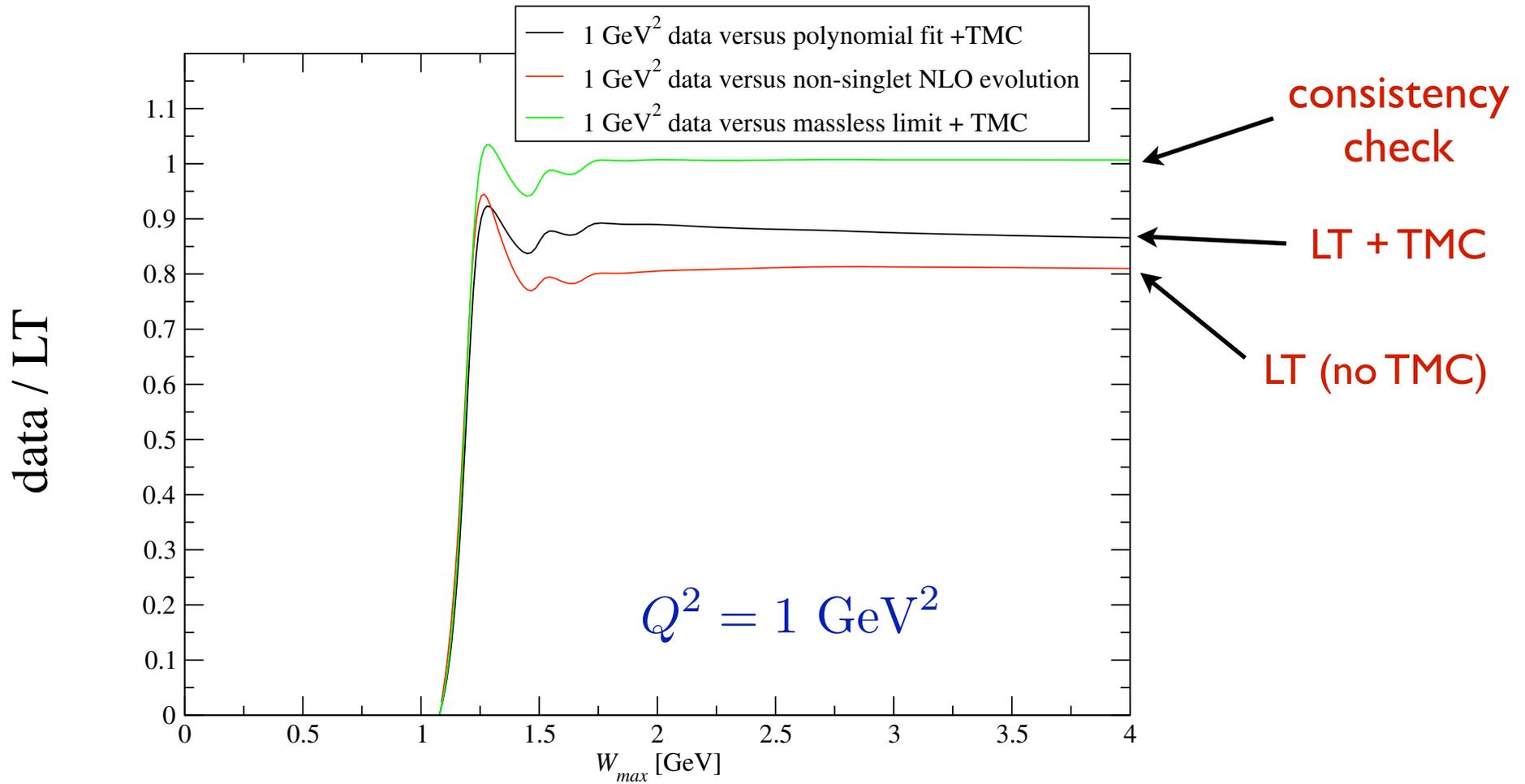
how much of this region
is leading twist ?

Analysis of JLab data

- assume data at highest Q^2 ($Q^2 = 9 \text{ GeV}^2$) is entirely leading twist
 - evolve (as NS) fit to data at $Q^2 = 9 \text{ GeV}^2$ down to lower Q^2
- apply TMC, and compare with data at lower Q^2



■ ratio of data to leading twist



- consider individual resonance regions:

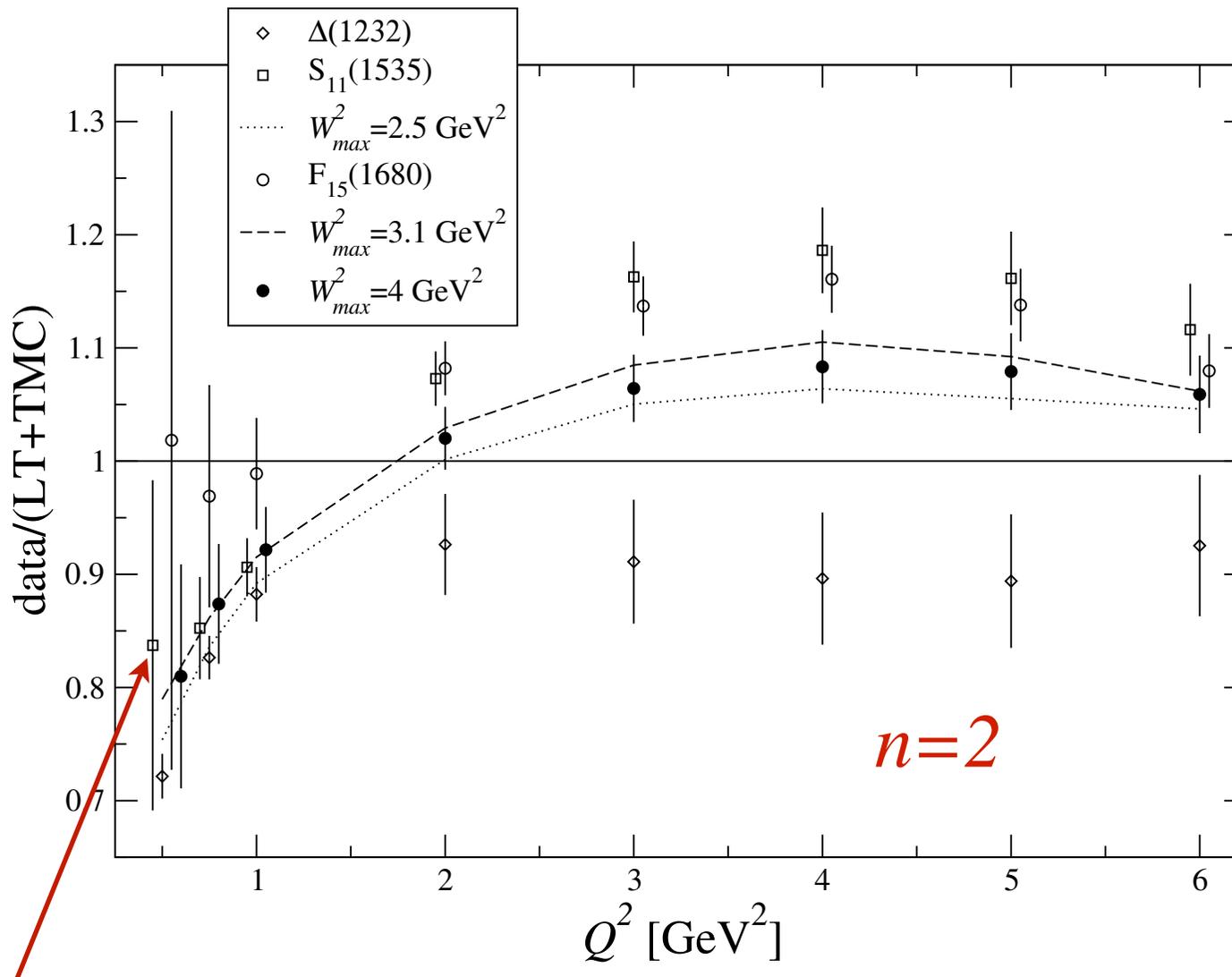
$$W_{\text{thr}}^2 < W^2 < 1.9 \text{ GeV}^2 \quad \text{“}\Delta(1232)\text{”}$$

$$1.9 < W^2 < 2.5 \text{ GeV}^2 \quad \text{“}S_{11}(1535)\text{”}$$

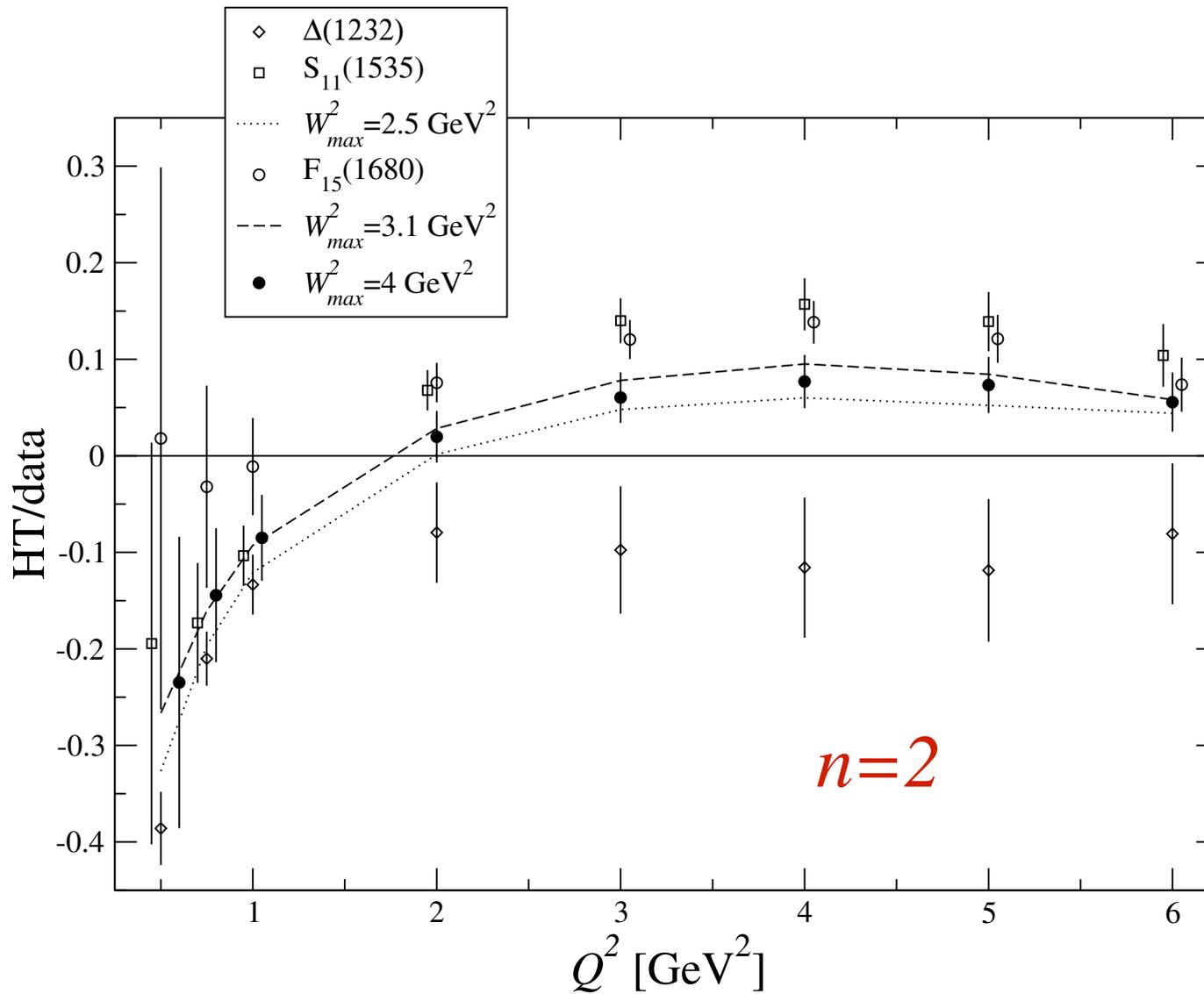
$$2.5 < W^2 < 3.1 \text{ GeV}^2 \quad \text{“}F_{15}(1680)\text{”}$$

as well as total resonance region:

$$W^2 < 4 \text{ GeV}^2$$

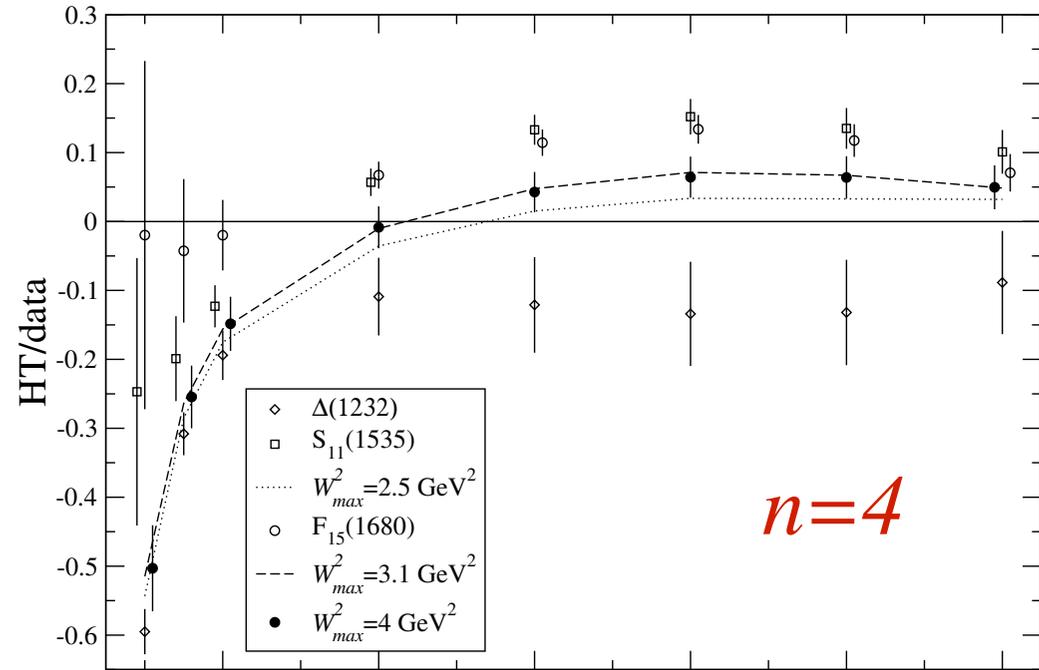


method breaks down for
 low x (high W) at low Q^2

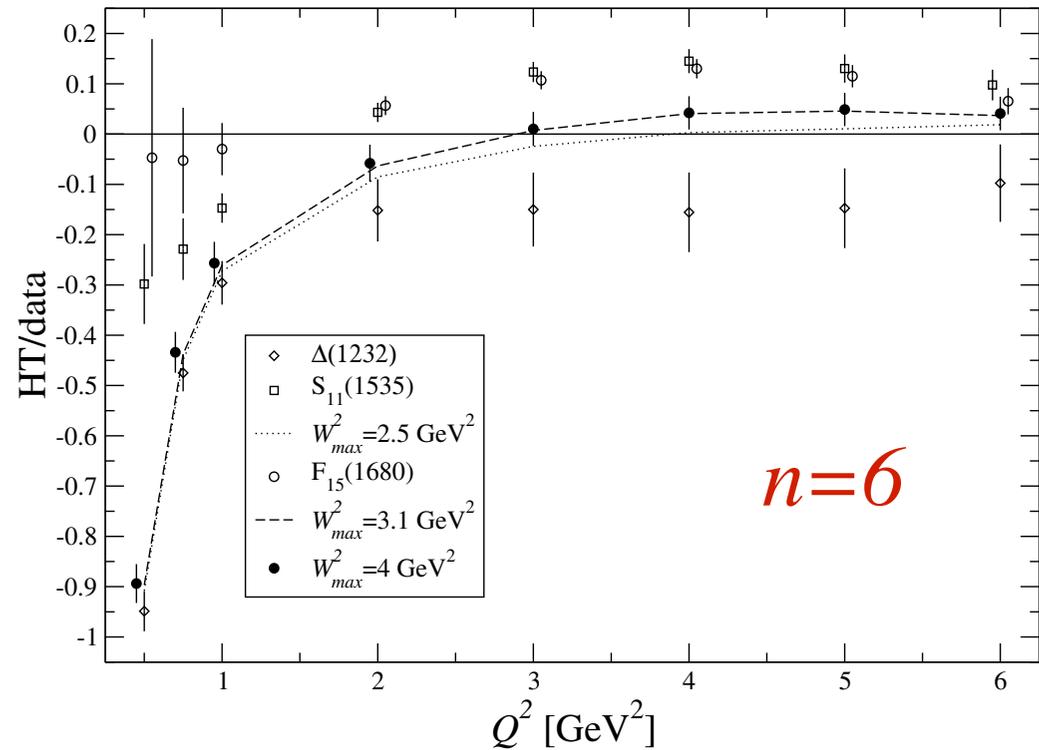


\rightarrow higher twists < 10% for $Q^2 > 1 \text{ GeV}^2$

■ higher moments



→ higher twists < 5-10%
for $Q^2 > 2-3$ GeV²
in resonance region



Summary

- Observation of quark-hadron duality in structure functions
 - higher twists “small” down to low Q^2 ($\sim 1 \text{ GeV}^2$)
 - global duality understood within QCD moments
- Local duality
 - duality exists in local regions of x (or W)
 - difficult to understand within QCD
- Truncated moments
 - firm foundation for study of local duality in QCD
 - higher twists $< 10\%$ for $Q^2 > 1 \text{ GeV}^2$ in resonance region